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## Molecular Structure and Intramolecular Electron Delocalization

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## 2.1 A Primer on Quantum Mechanics

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# Description of Quantum-Mechanical Systems

- **wave-particle duality** implies that **electromagnetic irradiation** can display **particle nature** while particles like electrons can display the characteristics of waves.
- all information on the quantum state of a quantum system at any time  $t$  can be described mathematically by a **complex function**  $\psi(r, t)$ , called **wavefunction**
- $\psi(r, t)$  can be regarded as a **probability amplitude** that defines the **probability of the possible results of measurements** made on the quantum system

- Dirac (“bra-ket”) notation treats wavefunctions as state vectors and utilizes bras and kets
  - ket  $|\phi\rangle$  denotes a state vector, that is, an element of an abstract complex vector space, the state space  $V$ , and represents a state of a quantum system
  - bra  $\langle f|$  denotes a linear functional  $f: V \rightarrow \mathbb{C}$  that maps each vector  $|\phi\rangle$  in the vector space  $V$  to a complex number in the complex plane  $\mathbb{C}$ , as given by the scalar product  $(\langle f|, |\phi\rangle)$ , denoted most often  $\langle f|\phi\rangle$
  - letting the linear functional  $\langle f|$  act on a vector  $|\nu\rangle$  is denoted as  $\langle f|\nu\rangle \in \mathbb{C}$  that has the form of a matrix multiplication of the row vector  $\langle f|$  with the column vector  $|\nu\rangle$
- advantages to understanding wavefunctions as elements of an abstract vector space:
  - Bra-ket notation accomplishes simpler formulations of wavefunctions
  - linear algebra can be used to manipulate and understand state vectors corresponding to wavefunctions

# State Space

- the ensemble of all wavefunctions, or quantum states, in which a quantum system can be found, forms an abstract complex vector space, the **state space**  $V$
- $V$  is a **Hilbert space**, i.e., a complex vector space equipped with a **scalar product** (inner product, projection product), that is a map  $V \times V \rightarrow \mathbb{C}$  defined as

$$(\psi_1, \psi_2) = \int_{-\infty}^{\infty} \psi_1^*(r, t) \psi_2(r, t) dr$$

- state spaces can be of finite dimension or infinite dimension; in the latter case, state spaces can also be discrete or continuous.
  - The spin state space of an electron in a magnetic field is of dimension 2, constituted of the state up and the state down.
  - The orbital state space of an electron in an infinite potential is infinite and discrete.
  - The position state space of an electron in space is infinite and continuous.

# Properties of the Scalar Product

## Inversion and Linearity

- $\langle \psi | \phi \rangle = \langle \phi | \psi \rangle^*$
- linearity on the right:  $\langle \psi | \lambda_1 \phi_1 + \lambda_2 \phi_2 \rangle = \lambda_1 \langle \psi | \phi_1 \rangle + \lambda_2 \langle \psi | \phi_2 \rangle$  where  $(\lambda_1, \lambda_2) \in \mathbb{C}^2$
- semi-linearity on the left:  $\langle \lambda_1 \psi_1 + \lambda_2 \psi_2 | \phi \rangle = \lambda_1^* \langle \psi_1 | \phi \rangle + \lambda_2^* \langle \psi_2 | \phi \rangle$  where  $(\lambda_1, \lambda_2) \in \mathbb{C}^2$

## Normalization

- $\langle \phi | \phi \rangle \in \mathbb{R}^+$ .  $\sqrt{\langle \phi | \phi \rangle}$  also noted  $\|\phi\|$  is called the norm of  $|\phi\rangle$
- a state  $|\phi\rangle$  is said to be normalised if  $\|\phi\| = 1$

## Orthogonality

- $\langle \phi | \phi \rangle = 0 \Leftrightarrow |\phi\rangle = 0$
- two states  $|\phi\rangle$  and  $|\psi\rangle$  are said to be orthogonal if  $\langle \psi | \phi \rangle = 0$

# Observables and Expectation Values

- any **observable  $\Omega$** , that is, a **mesurable physical property of a quantum system**, can be represented by a corresponding **operator  $\hat{\Omega}$** , a mathematical operation that can be performed on the wave function  $\psi$  respectively the corresponding state vector  $|\psi\rangle$
- to obtain more information on an observable  $\hat{\Omega}$  of a defined quantum system, the following associated equation, called **Eigenvalue equation**, needs to be solved:

$$\hat{\Omega} |\psi\rangle = \omega |\psi\rangle$$

- wavefunctions  $\psi$  that **fulfill the Eigenvalue equation** are called **Eigenfunctions** (allowed states)
- **Eigenvalue  $\omega$**  is a constant that to the **value of the observable  $\Omega$**
- **expectation value** can be calculated as

$$\langle \Omega \rangle = \int \Psi^* \hat{\Omega} \Psi d\tau = \omega$$

because if  $\Psi$  is an **Eigenfunction**, then  $\int \Psi^* \omega \Psi d\tau = \omega \int \Psi^* \Psi d\tau = \omega$

# Hermitian Conjugate Operators and Hermiticity

- for any linear operator  $\hat{\Omega}$ , the conjugate operator (adjoint)  $\hat{\Omega}^\dagger$  is defined as:

$$\langle \Psi_i | \hat{\Omega} \Psi_j \rangle = \int \Psi_i^* \hat{\Omega} \Psi_j d\tau = \int \Psi_j (\hat{\Omega}^\dagger \Psi_i)^* d\tau = \langle \hat{\Omega}^\dagger \Psi_i | \Psi_j \rangle$$

- **Hermiticity:** quantum-mechanical operators  $\hat{\Omega}$  that are **identical to their own adjoint** are said to be **Hermitian operators**.

$$\hat{\Omega} = \hat{\Omega}^\dagger$$

- **any quantum-mechanical operator  $\hat{\Omega}$  corresponding to a physical observable  $\Omega$  can be shown to be a Hermitian operator.**

# Consequences of Hermiticity

- Hermiticity has important implications (by its definition):
  - all **Eigenvalues**, the values of the observables, are **always real**, that is,  $\omega = \omega^*$
  - the set of all **Eigenfunctions**  $\Psi(r, t)$  is **complete**
  - the ensemble of the Eigenfunctions form an **orthonormal basis of the state**, that is, they are orthogonal (hence, linearly independent) and normalized

$$\langle \Psi_i | \Psi_j \rangle = \int \Psi_i^* \Psi_j d\tau = 0 \quad \text{for } i \neq j \quad \text{and} \quad \langle \Psi_i | \Psi_i \rangle = \int \Psi_i^* \Psi_i d\tau = 1$$

- any state vector can be formed from a combination of basis states with complex coefficients:

$$|\phi\rangle = \sum_{\nu} c_{\nu} |\psi_{\nu}\rangle \quad \text{with } c_{\nu} \in \mathbb{C}, \text{ for a discrete case}$$

$$|\phi\rangle = \int c(r) |r\rangle dr \quad \text{with } c(r) \in \mathbb{C}, \text{ for a continuous case}$$

- **Hermiticity of quantum-mechanical operators  $\hat{\Omega}$  corresponding to a physical observables  $\Omega$  are the basis for concepts such as hybridization and linear combination of atomic orbitals**

# Born Interpretation of the Wavefunction

- probability density  $\rho(r, t)$  to find a particle at the location  $r$  at time  $t$  is given by the square modulus of the complex wave function (which implies the density is always a real number)

$$\rho(r, t) = |\psi(r, t)|^2 = \psi^*(r, t)\psi(r, t)$$

- probability  $P(t)$  to find a particle in an infinitesimal volume  $d\tau$  around a location  $r$  at time  $t$  is

$$P(t) = \int |\psi|^2 d\tau$$

- normalization: the probability integrated over the entire space equals to 1

$$\int |N\psi|^2 d\tau = 1 \Leftrightarrow N = \left\{ \int |N\psi|^2 d\tau \right\}^{-1/2}$$

- due to Hermiticity, Eigenfunctions resulting from Eigenequations with quantum-mechanical operators  $\hat{\Omega}$  corresponding to a physical observables  $\Omega$  are normalized

# Quantization Resulting from the Constraints of the Born Interpretation

- an acceptable wavefunction must hence be square-integrable and normalizable function mapping each point of 3D space to a complex number
- Born interpretation further implies that the wavefunction is
  - single-valued at every location  $r$
  - not infinite over a non-infinitesimal region
  - continuous in slope and curvature
- **these conditions limits the choice of acceptable mathematical functions**
- **quantization of the allowed energies of a particle is the result of the finite probability density and the resulting constraints for the selection of acceptable wavefunctions**